Adaptive finite elements with high aspect ratio for the computation of dendritic growth and coalescence in binary alloys.

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Outline

- A phase field model for dendritic growth of binary alloys
 - ★ Model
 - ★ Existence, a priori and a posteriori error estimates
 - ★ Numerical Tool : Adaptive finite elements with high aspect ratio
- A multiphase field model for the coalescence of binary alloys

Solidification : dendritic growth



Sa-Co alloy

Solidification : from macro to micro scale



Grain shape and size (\rightarrow mechanical properties) : depend on solidification history

A phase field model for dendritic growth of a binary alloy

• Phase field model (Warren Boettinger Acta Metall. Mater. 1995, Tiaden Nestler Diepers Steinbach Phys D 1998). Find $c, \phi : \Omega \times (0, T) \to \mathbb{R}$ such that

$$\frac{\partial \phi}{\partial t} - \operatorname{div} \left(A(\nabla \phi) \nabla \phi \right) \right) - S(c, \phi) = 0 \qquad \text{in } \Omega \times (0, T),$$
$$\frac{\partial c}{\partial t} - \operatorname{div} \left(D_1(\phi) \nabla c + D_2(c, \phi) \nabla \phi \right) = 0 \qquad \text{in } \Omega \times (0, T).$$



• interface width : ε small $\simeq 10^{-8} m$ (100 atoms).

The phase field equation

• Phase field

$$\frac{\partial \phi}{\partial t} - \operatorname{div} \left(A(\nabla \phi) \nabla \phi) \right) \ - S(c, \phi) = 0.$$

• Free energy = interface (anisotropy) + double wells + phase transformation

$$\frac{\varepsilon}{2} \int_{\Omega} a^2(\theta(\nabla\phi)) |\nabla\phi|^2 + \frac{1}{\varepsilon} \int_{\Omega} W(\phi) + \int_{\Omega} c_\ell - c_\ell^{eq},$$

with $a(\theta) = (1 + \bar{a}\cos(\kappa\theta))$ and $\theta(\nabla\phi)$ angle between $\nabla\phi$ and horizontal.

- Small parameter in $S(c, \phi)$ (time step $\simeq \varepsilon^2$, Elliott State of the art in numer. anal. 1997), strongly nonlinear div $(A(\nabla \phi)\nabla \phi))$.
- Existence for low physical anisotropy $\bar{a} < \frac{1}{\kappa^2 1}$, a priori error estimates (Burman Rappaz M3AS 2003).
- A posteriori error estimates, adaptive FE with high aspect ratio (Burman Picasso J. Interfaces Free Boundaries 2003).

Small and large physical anisotropy

- Interface energy $\int_{\Omega} a^2(\theta(\nabla \phi)) |\nabla \phi|^2$.
- $a(\theta) = (1 + \bar{a}\cos(\kappa\theta))$ and $\theta(\nabla\phi)$ angle between $\nabla\phi$ and horizontal.
- isolines of $\xi \to a^2(\theta(\xi)) |\xi|^2$





Small physical anisotropy $\bar{a} < \frac{1}{\kappa^2 - 1}$

Large physical anisotropy $\bar{a} > \frac{1}{\kappa^2 - 1}$

A posteriori error estimates, adaptive FE with high aspect ratio

- A posteriori error estimates : error \leq estimator(mesh,numerical solution).
- $A(\cdot)$ is strongly elliptic : $\bar{a} < \frac{1}{\kappa^2 1}$, $\exists \mu$

$$\mu \int_{\Omega} |\nabla(\phi - \psi)|^2 \le \int_{\Omega} \Big(A(\nabla\phi) \nabla\phi - A(\nabla\psi) \nabla\psi \Big) \cdot \nabla(\phi - \psi) \qquad \forall \phi, \psi \in H^1_0(\Omega)$$

- The error in the $L^2(0,T;L^2(\Omega))$ norm converges faster to zero than the error in the $L^2(0,T;H^1(\Omega))$ norm.
- Then, $\exists C$ independent of "everything" (problem data, mesh size and aspect ratio) such that

$$\begin{split} \mu \int_0^T \int_\Omega |\nabla(\phi - \phi_h)|^2 + \frac{\mu D_1^2}{D_2^2} \int_0^T \int_\Omega |\nabla(c - c_h)|^2 \\ & \leq C \left(est_\phi(mesh, \phi_h, c_h) + \frac{\mu D_1^2}{D_2^2} est_c(mesh, \phi_h, c_h) \right). \end{split}$$

Error indicator for meshes with high aspect ratio



$$\frac{1}{2\lambda_{2,K}^{1/2}} \left\| \left[\frac{\partial c_h}{\partial n} \right] \right\|_{L^2(\partial K)} \left(\lambda_{1,K}^2 \left(\mathbf{r}_{1,K}^T G_K(c-c_h) \mathbf{r}_{1,K} \right) + \lambda_{2,K}^2 \left(\mathbf{r}_{2,K}^T G_K(c-c_h) \mathbf{r}_{2,K} \right) \right)^{1/2} \right\|_{L^2(\partial K)} \left(\lambda_{1,K}^2 \left(\mathbf{r}_{1,K}^T G_K(c-c_h) \mathbf{r}_{1,K} \right) + \lambda_{2,K}^2 \left(\mathbf{r}_{2,K}^T G_K(c-c_h) \mathbf{r}_{2,K} \right) \right)^{1/2} \right\|_{L^2(\partial K)} \left\| \left(\lambda_{1,K}^2 \left(\mathbf{r}_{1,K}^T G_K(c-c_h) \mathbf{r}_{1,K} \right) + \lambda_{2,K}^2 \left(\mathbf{r}_{2,K}^T G_K(c-c_h) \mathbf{r}_{2,K} \right) \right)^{1/2} \right\|_{L^2(\partial K)} \left\| \left(\lambda_{1,K}^2 \left(\mathbf{r}_{1,K}^T G_K(c-c_h) \mathbf{r}_{1,K} \right) + \lambda_{2,K}^2 \left(\mathbf{r}_{2,K}^T G_K(c-c_h) \mathbf{r}_{2,K} \right) \right)^{1/2} \right\|_{L^2(\partial K)} \right\|_{L^2(\partial K)} \left\| \left(\lambda_{1,K}^2 \left(\mathbf{r}_{1,K}^T G_K(c-c_h) \mathbf{r}_{1,K} \right) + \lambda_{2,K}^2 \left(\mathbf{r}_{2,K}^T G_K(c-c_h) \mathbf{r}_{2,K} \right) \right)^{1/2} \right\|_{L^2(\partial K)} \right\|_{L^2(\partial K)} \left\| \left(\lambda_{1,K}^2 \left(\mathbf{r}_{1,K}^T G_K(c-c_h) \mathbf{r}_{1,K} \right) + \lambda_{2,K}^2 \left(\mathbf{r}_{2,K}^T G_K(c-c_h) \mathbf{r}_{2,K} \right) \right\|_{L^2(\partial K)} \right\|_{L^2(\partial K)} \right\|_{L^2(\partial K)} \left\| \left(\lambda_{1,K}^2 \left(\mathbf{r}_{1,K}^T G_K(c-c_h) \mathbf{r}_{1,K} \right) + \lambda_{2,K}^2 \left(\mathbf{r}_{2,K}^T G_K(c-c_h) \mathbf{r}_{2,K} \right) \right\|_{L^2(\partial K)} \right\|_{L^2(\partial K)} \right\|_{L^2(\partial K)} \left\| \left(\lambda_{1,K}^2 \left(\mathbf{r}_{1,K}^T G_K(c-c_h) \mathbf{r}_{1,K} \right) + \lambda_{2,K}^2 \left(\mathbf{r}_{2,K}^T G_K(c-c_h) \mathbf{r}_{2,K} \right) \right\|_{L^2(\partial K)} \right\|_{L^2(\partial K)} \left\| \left(\lambda_{1,K}^2 \left(\mathbf{r}_{1,K}^T G_K(c-c_h) \mathbf{r}_{2,K} \right) \right) \right\|_{L^2(\partial K)} \right\|_{L^2(\partial K)} \right\|_{L^2(\partial K)} \left\| \left(\lambda_{1,K}^2 \left(\mathbf{r}_{1,K}^T G_K(c-c_h) \mathbf{r}_{2,K} \right) \right) \right\|_{L^2(\partial K)} \right\|_{L^2(\partial K)} \right\|_{L^2(\partial K)} \left\| \left(\lambda_{1,K}^2 \left(\mathbf{r}_{1,K}^T G_K(c-c_h) \mathbf{r}_{2,K} \right) \right\|_{L^2(\partial K)} \right\|_{L^2(\partial K)} \right\|_{L^2(\partial K)} \left\| \left(\lambda_{1,K}^2 \left(\mathbf{r}_{1,K}^T G_K(c-c_h) \mathbf{r}_{2,K} \right) \right\|_{L^2(\partial K)} \right\|_{L^2(\partial K)} \right\|_{L^2(\partial K)} \left\| \left(\lambda_{1,K}^2 \left(\mathbf{r}_{1,K}^2 G_K(c-c_h) \mathbf{r}_{2,K} \right) \right\|_{L^2(\partial K)} \right\|_{L^2(\partial K)} \right\|_{L^2(\partial K)} \left\| \left(\lambda_{1,K}^2 \left(\mathbf{r}_{1,K}^2 G_K(c-c_h) \mathbf{r}_{2,K} \right) \right\|_{L^2(\partial K)} \right\|_{L^2(\partial K)} \right\|_{L^2(\partial K)} \right\|_{L^2(\partial K)} \right\|_{L^2(\partial K)} \left\| \left(\lambda_{1,K}^2 G_K(c-c_h) \right\|_{L^2(\partial K)} \right\|_{L^2(\partial K)} \right\|_{L^2(\partial K)} \left\| \left(\lambda_{1,K}^2 G_K(c-c_h) \mathbf{r}_{2,K} \right) \right\|_{L^2(\partial K)} \right\|_{L^2(\partial K)} \right\|_{L^2(\partial K)} \right\|_{L^2(\partial K)} \left\| \left\| \left(\lambda_{1,K}^2 G_K(c-c_h) \right\|_{L^2(\partial K)} \right\|_{L^2(\partial K)} \right\|_{L$$

where $G_K(\cdot)$ is the matrix of first order derivatives, approached using post-processing techniques (Picasso SISC 2003).

Numerical validation of the error indicator



ei : effectivity index, error indicator / true error.

Small physical anisotropy

• Goal of the adaptive algorithm : build a sequence of triangulations such that

 $\frac{\text{error indicator}}{\int_0^T \int_\Omega |\nabla c_h|^2} \simeq TOL.$

• Box width 5 10^{-4} , $\varepsilon = 5 \ 10^{-7}$, time step 5 10^{-4} , TOL = 0.0625, 24000 vertices, final time 1.



Small physical anisotropy





Same precision with isotropic meshes : 10 times more vertices !

Convergence





TOL=0.25 1987 Vertices TOL=0.125 6073 Vertices



TOL=0.0625 24441 Vertices



TOL=0.03125 101782 Vertices

Large physical anisotropy





ϕ isolines according to physical anisotropy



A multiphase field model for coalescence in binary alloys



- Multiphase field : Tiaden Nestler Diepers Steinbach Phys D 1998
- Multiphase field for binary alloys : Rappaz Jacot Boettinger Met. Trans. A 2003, Burman Jacot Picasso J. Comp. Phys. 2004
- Unknowns : ϕ_1 , ϕ_2 , ϕ_3 , λ (Lagrange multiplier $\phi_1 + \phi_2 + \phi_3 = 1$) and c.

Free energy

• Interface energy

$$J_{int}(\phi_1, \phi_2, \phi_3) = \frac{\varepsilon_{12}^2}{2} \int_{\Omega} a^2 \Big(\theta_{12} \big(\mathbf{r}(\phi_1, \phi_2) \big) \Big) |\mathbf{r}(\phi_1, \phi_2)|^2 \\ + \frac{\varepsilon_{13}^2}{2} \int_{\Omega} a^2 \Big(\theta_{13} \big(\mathbf{r}(\phi_1, \phi_3) \big) \Big) |\mathbf{r}(\phi_1, \phi_3)|^2 \\ + \frac{\varepsilon_{23}^2}{2} \int_{\Omega} a^2 \Big(\theta_{23} \big(\mathbf{r}(\phi_2, \phi_3) \big) \Big) |\mathbf{r}(\phi_2, \phi_3)|^2$$

where $\mathbf{r}(\phi_i, \phi_j) = \phi_i \nabla \phi_j - \phi_j \nabla \phi_i$ $(i, j) \in \{(1, 2), (1, 3), (2, 3)\}$

• Phase transformation $J_{tr}(\phi_1, \phi_2, \phi_3) = \int_{\Omega} (c_{\ell} - c_{\ell}^{eq})$

• Double wells
$$J_{dw}(\phi_1, \phi_2, \phi_3) = \frac{1}{\varepsilon} \int_{\Omega} \phi_1^2 \phi_2^2 + \phi_1^2 \phi_3^2 + \phi_2^2 \phi_3^2$$

The parabolic system of equations

- Find ϕ_1 , ϕ_2 , ϕ_3 such that $J_{int} + J_{tr} + J_{dw}$ =min and $\phi_1 + \phi_2 + \phi_3 = 1$
- Lagrangian

$$\mathcal{L}(\phi_1, \phi_2, \phi_3, \lambda) = (J_{int} + J_{tr} + J_{dw})(\phi_1, \phi_2, \phi_3) + \int_{\Omega} \lambda(\phi_1 + \phi_2 + \phi_3 - 1).$$

- $\mathcal{DL}(\phi_1, \phi_2, \phi_3, \lambda) = 0$
- Find ϕ_1 , ϕ_2 , ϕ_3 , λ such that

$$\begin{aligned} \frac{\partial \phi_1}{\partial t} + \mathcal{D}_{\phi_1} \mathcal{L}(\phi_1, \phi_2, \phi_3, \lambda) &= 0, \\ \frac{\partial \phi_2}{\partial t} + \mathcal{D}_{\phi_2} \mathcal{L}(\phi_1, \phi_2, \phi_3, \lambda) &= 0, \\ \frac{\partial \phi_3}{\partial t} + \mathcal{D}_{\phi_3} \mathcal{L}(\phi_1, \phi_2, \phi_3, \lambda) &= 0, \\ \phi_1 + \phi_2 + \phi_3 - 1 &= 0. \end{aligned}$$

The parabolic system of equations

• Weak formulation with no anisotropy

$$\begin{split} \int_{\Omega} & \left(\begin{pmatrix} \frac{\partial \phi_1}{\partial t} \\ \frac{\partial \phi_2}{\partial t} \\ \frac{\partial \phi_3}{\partial t} \end{pmatrix} \cdot \begin{pmatrix} \psi_1 \\ \psi_2 \\ \psi_3 \end{pmatrix} + \begin{pmatrix} (\phi_2^2 + \phi_3^2)I_2 & -\phi_1\phi_2I_2 & -\phi_1\phi_3I_2 \\ & \dots & & \end{pmatrix} \begin{pmatrix} \nabla \phi_1 \\ \nabla \phi_2 \\ \nabla \phi_3 \end{pmatrix} \cdot \begin{pmatrix} \nabla \psi_1 \\ \nabla \psi_2 \\ \nabla \psi_3 \end{pmatrix} \\ & + \begin{pmatrix} |\nabla \phi_2|^2 + |\nabla \phi_3|^2 & -\nabla \phi_1 \cdot \nabla \phi_2 & -\nabla \phi_1 \cdot \nabla \phi_3 \\ & \dots & & \end{pmatrix} \begin{pmatrix} \phi_1 \\ \phi_2 \\ \phi_3 \end{pmatrix} \cdot \begin{pmatrix} \psi_1 \\ \psi_2 \\ \psi_3 \end{pmatrix} \\ & + \text{double well stuff} + \lambda(\psi_1 + \psi_2 + \psi_3) \end{pmatrix} = 0, \end{split}$$

for all $\overline{\psi_1,\psi_2,\psi_3}$.

$$\int_{\Omega} (\phi_1 + \phi_2 + \phi_3)\mu = \int_{\Omega} \mu,$$

for all μ .



Box $3\ 10^{-5}$, $\varepsilon = 5\ 10^{-8}$, time step 10^{-3} , TOL = 0.25, 3000 vertices time 2.

Results





Results



Conclusions and perspectives

- Dendritic growth and coalescence : adaptive finite elements with high aspect ratio \rightarrow accurate results with few vertices (< 20 000).
- Mathematical analysis of coalescence.
- Dendritic growth and coalescence with convection.
- 3D ?

Dendritic growth and coalescence with convection



Concentration



Dendritic growth and coalescence with convection



Concentration

